

# Simulations of Low- $q$ Disruptions in the Compact Toroidal Hybrid Experiment

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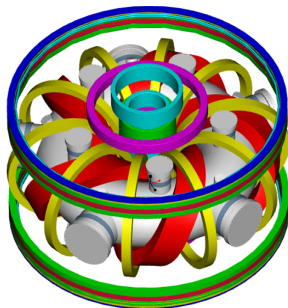
# Numerical simulations are used to investigate the MHD precursors to low- $q$ disruptions in the Compact Toroidal Hybrid device (CTH).

- Experimental Observations
- Numerical Model
- Simulation Results
  - Disrupting  $t_{vac} = 0.015$  Discharge
  - Non-Disrupting  $t_{vac} = 0.075$  Discharge
- Conclusions

The Compact Toroidal Hybrid device (CTH) is stellarator-tokamak hybrid designed to study the effects of 3D shaping on MHD instabilities.

CTH Parameters

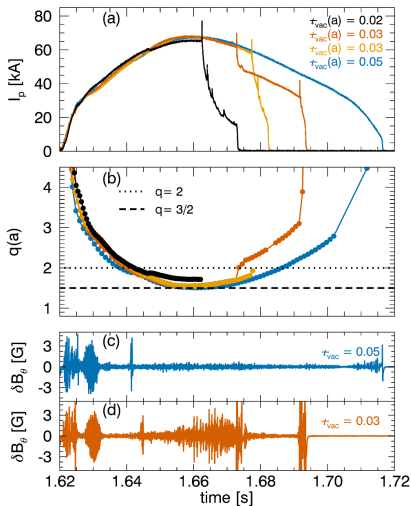
Field Periods	5
Major Radius	0.75m
Minor Radius	0.20m
Magnetic Field	$\leq 0.7$ T
Plasma Current	$\leq 80$ kA
Number Density	$\leq 5 \times 10^{19} \text{m}^{-3}$
Electron Temperature	$\leq 200$ eV



- The rotational transform is generated by a combination of external 3D helical coil currents and internal plasma currents.
  - The rotational transform,  $t$ , is the inverse of the safety factor:  $t = 1/q$ .
- The fractional transform,  $f$ , quantifies the amount of 3D shaping.
  - $f = t_{vac}/t_{total}$
- CTH can operate with a fractional transform that ranges from  $f = 4\%$  to  $f = 100\%$  by adjusting the plasma current.

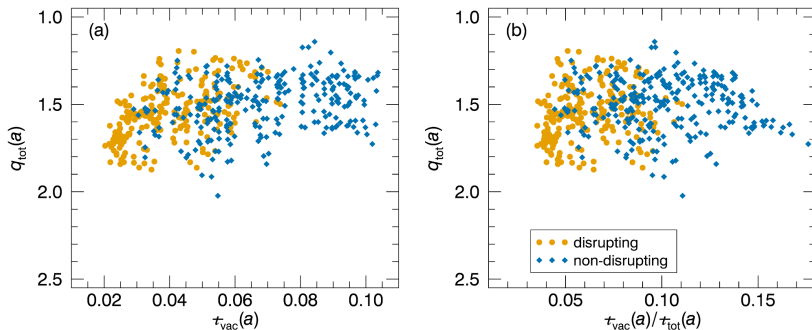
# Small amounts of externally applied rotational transform allows CTH to operate with $q(a) < 2$ .

- External kink stability limits tokamak operation to  $q(a) \geq 2$ .
- Disruptions are observed in low- $q$  discharges after peak plasma current.
  - 3/2 mode activity is observed in both disrupting and non-disrupting discharges.
  - 4/3 mode activity is only observed in disrupting discharges.
  - 1/1 activity is observed in both cases.
- Disruptions occur when the edge safety factor passes through  $q(a) \approx 1.7$ .



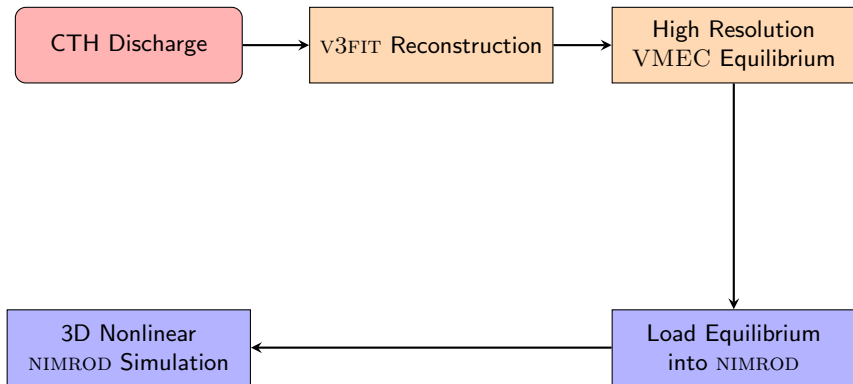
[M. D. Pandya et al., POP 22, 2015 ]

# Low- $q$ disruptions are suppressed at large vacuum transform.



- The frequency of disruptions decreases with increasing vacuum transform.
  - Disruptions always occur when  $\tau_{vac} \lesssim 0.03$ .
  - Disruptions are completely suppressed for  $\tau_{vac} \gtrsim 0.07$ .
- Here  $q_{tot}(a)$  is the value of the edge safety factor at peak plasma current.

Nonlinear simulations are initialized with  $v3FIT$  reconstructions of experimental discharges.



$$\begin{aligned} \frac{\partial \vec{B}}{\partial t} &= -\nabla \times \left[ \eta \left( \vec{J} - \vec{J}_S \right) - \vec{V} \times \vec{B} \right] + k_{divb} \nabla \nabla \cdot \vec{B} \\ \frac{\partial n}{\partial t} + \nabla \cdot (n \vec{V}) &= \nabla \cdot (D \nabla n - D_h \nabla \nabla^2 n) \\ \rho \left( \frac{\partial \vec{V}}{\partial t} + \vec{V} \cdot \nabla \vec{V} \right) &= \vec{J} \times \vec{B} - \nabla P - \nabla \cdot \vec{\pi} \\ \frac{3}{2} n \left( \frac{\partial T}{\partial t} + \vec{V} \cdot \nabla T \right) &= -P \nabla \cdot \vec{V} - \nabla \cdot \vec{q} + \eta J^2 \end{aligned}$$

- 3D current source maintains the current profile against resistive decay.
- Magnetic divergence diffusivity enforces  $\nabla \cdot \vec{B} \approx 0$ .
- Artificial particle diffusivity smooths small scale density fluctuations.

# Transport effects are modeled using a collisional closure for a hydrogen plasma.

- Temperature dependent magnetic diffusivity:

- $\frac{\eta}{\mu_0} = \frac{\eta_0}{\mu_0} T^{-\frac{3}{2}}$

- $\frac{\eta_0}{\mu_0} = 1.23 \times 10^3 \frac{m^2}{s}$

- Thermal conduction is anisotropic with temperature dependent thermal diffusivities:

- $\vec{q} = -n\chi_{\parallel} \nabla_{\parallel} T - n\chi_{\perp} \nabla_{\perp} T$

- $\chi_{\parallel} = \chi_{\parallel 0} T^{5/2}, \quad \chi_{\parallel 0} = 1.26 \times 10^3 \frac{m^2}{s}$

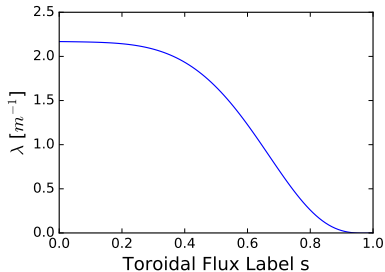
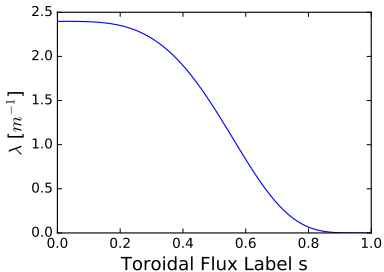
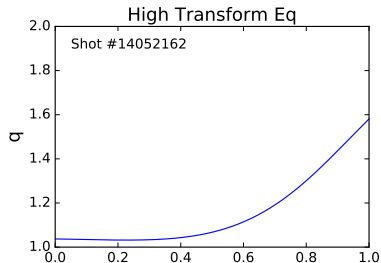
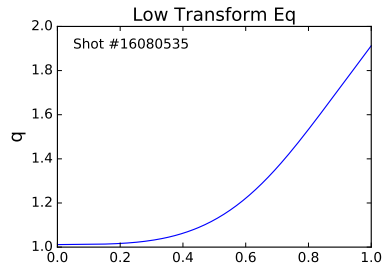
- $\chi_{\perp} = \chi_{\perp 0} T^{-1/2}, \quad \chi_{\perp 0} = 2.69 \frac{m^2}{s}$

- The viscous stress tensor is isotropic with a constant viscosity:

- $\vec{\pi} = -\nu_0 mn \left( \nabla \vec{V} + \nabla \vec{V}^T - \frac{2}{3} \vec{I} \nabla \cdot \vec{V} \right), \quad \nu_0 = 1 \frac{m^2}{s}$

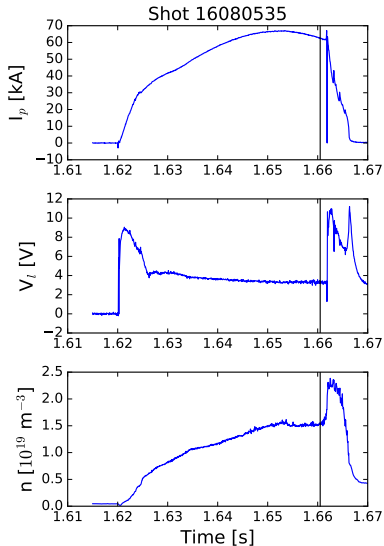


# Simulations of a small transform disrupting discharge are compared against simulations of a large transform non-disrupting discharge.

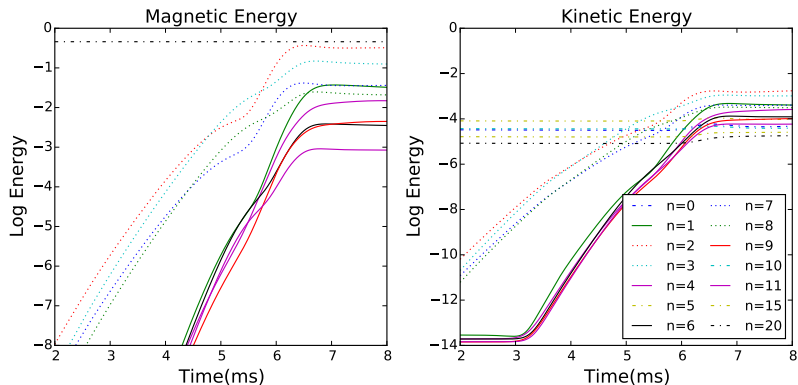


# A small $t_{vac}$ case is used to study the dynamics of a disrupting discharge.

- Simulations model a discharge with a  $t_{vac} = 0.015$ .
- Simulations are initialized with equilibrium reconstructions of the plasma 2ms before the disruption.
- Small  $t_{vac}$  discharges are most likely to disrupt.
- Small  $t_{vac}$  eases toroidal resolution requirements.

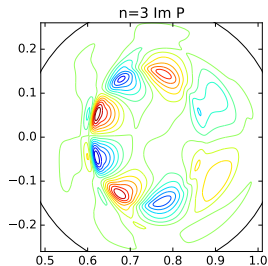
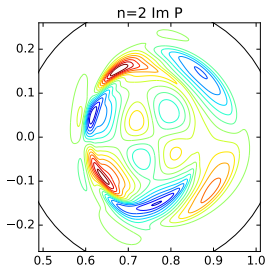
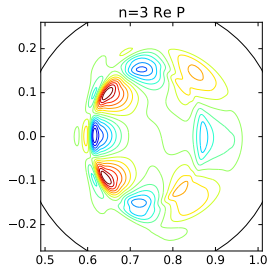
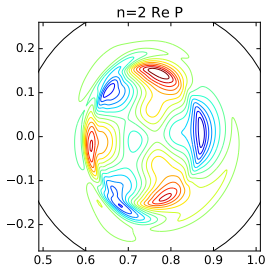


Predominantly  $n = 2$  and  $n = 3$  modes are unstable and saturate at large amplitude.

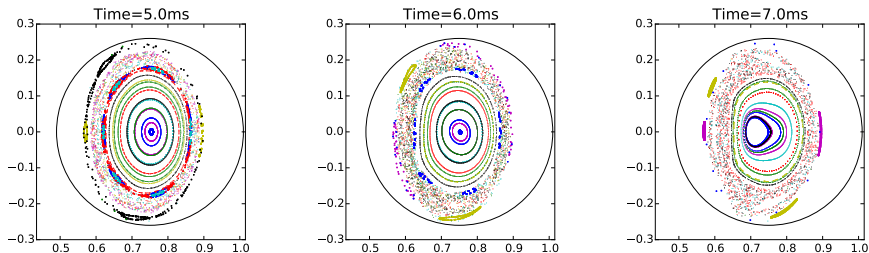


- Linear modes are composed of a family Fourier modes due to coupling with the stellarator equilibrium fields.

The dominant  $n = 2$  and  $n = 3$  poloidal structures are consistent with experimental observations.



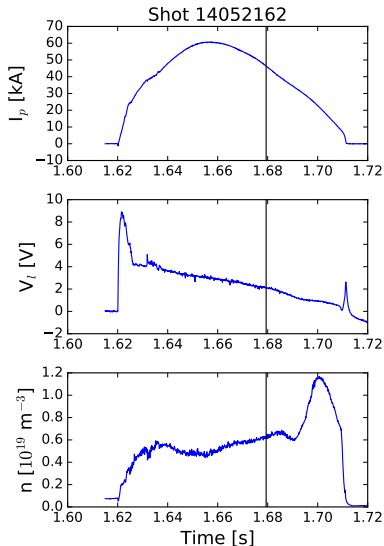
The interaction between the unstable modes and symmetry preserving islands leads to a large region of stochastic field.



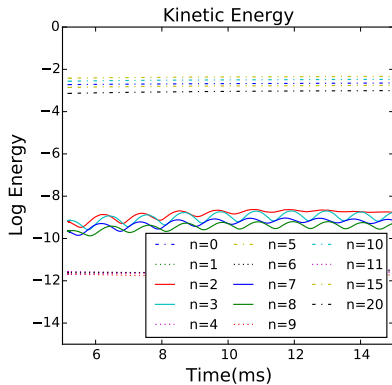
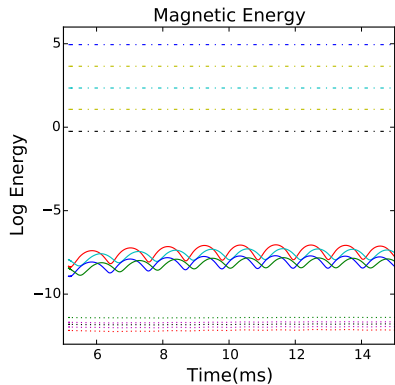
- Symmetry preserving islands are observed early in time.
  - 6/5, 7/5, 8/5, and 9/5 islands are observed
- Islands degrade as the unstable modes grow.
- Thermal energy is lost when the innermost island chain is destroyed.

# A large $t_{vac}$ discharge is used to study the dynamics of a non-disrupting discharge.

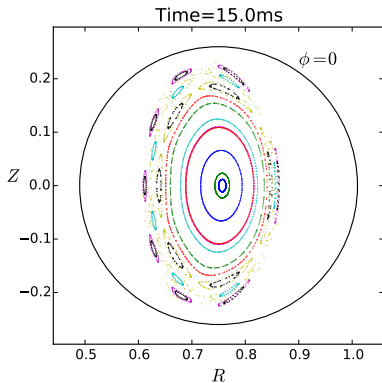
- This non-disrupting discharge has  $t_{vac} = 0.075$ .
- Disruptions are not observed in CTH discharges with  $t_{vac} \gtrsim 0.07$ .
- Simulations are initialized with reconstructions at conditions when  $q(a) \approx 1.7$ .
- In the experiment disruption occur most frequently when  $q(a) \approx 1.7$ .



The linearly unstable modes saturate at very small amplitude.



- The 6/5, 7/5, and 8/5 island chains are seen in the figure.
- Good nested flux surfaces are present in the core.
- Thermal energy increases throughout the simulation.





## Simulations suggest that low- $q$ disruptions are triggered by an interaction between symmetry preserving islands and symmetry breaking instabilities.

- Symmetry preserving islands form at the  $q = 6/5, 7/5, 8/5,$  and  $9/5$  rational surfaces.
- At small vacuum transform the  $4/3$  and  $3/2$  modes are unstable and grow to large amplitude.
  - The  $4/3$  and  $3/2$  modes interact with the symmetry preserving islands.
  - The destruction of the islands creates a large volume of stochastic field.
  - Thermal energy is lost when the inner most island chain is destroyed.
- At large vacuum transform the  $4/3$  and  $3/2$  modes saturate at small amplitude.
  - Symmetry preserving islands persist throughout the simulation.